

NIMROD simulations with simple current sources

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in collaboration with

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About NIMROD

- Evolves fluid equations for magnetically confined plasmas — stiff, anisotropic system
- Finite elements in poloidal plane; Fourier decomposition in toroidal direction
- Can be tailored to many plasma configurations: tokamak, spheromak, astrophysical, etc.

$$\begin{aligned}\frac{\partial \mathbf{B}}{\partial t} &= -\nabla \times \mathbf{E} + \kappa_b \nabla(\nabla \cdot \mathbf{B}) & \mathbf{E} + \mathbf{v} \times \mathbf{B} &= \eta \mathbf{J} \\ \frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{v}) &= \nabla \cdot (D\nabla n) & \mu_0 \mathbf{J} &= \nabla \times \mathbf{B} \\ \rho \frac{\partial \mathbf{v}}{\partial t} + \rho(\mathbf{v} \cdot \nabla)\mathbf{v} &= -\nabla p + \mathbf{J} \times \mathbf{B} + \nabla(\nu \rho \nabla \mathbf{v}) \\ \frac{n}{\gamma - 1} \left(\frac{\partial T}{\partial t} + (\mathbf{v} \cdot \nabla)T \right) &= -\frac{p}{2} \nabla \cdot \mathbf{v} - \nabla \cdot \mathbf{q} + Q\end{aligned}$$

- Numerical dissipation introduced

Insert an ad hoc term (general RF effects) in Ohm's law

- MHD: Ohm's law:

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{J}$$

- Determine effect of current drive (not self-consistent): ad hoc source term

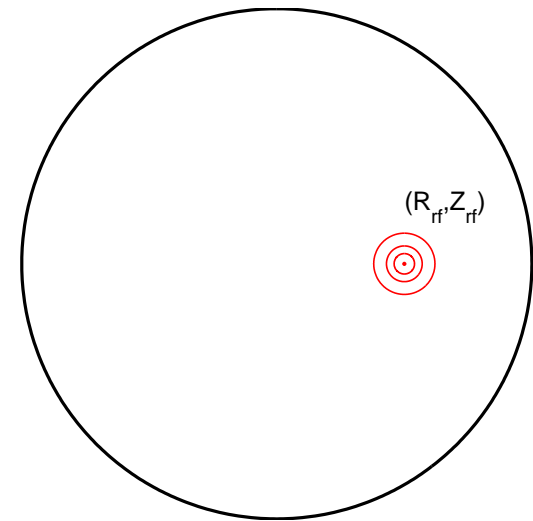
$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta(\mathbf{J} - \mathbf{J}_{\text{rf}})$$

- Should modify plasma equilibrium. Assume axisymmetric current profile of form

$$\mathbf{J}_{\text{rf}}(R, Z, t) = \lambda_{\text{rf}} \exp \left[\frac{- [(R - R_{\text{rf}})^2 + (Z - Z_{\text{rf}})^2]}{w_{\text{rf}}^2} \right] \frac{\mathbf{B}}{\mu_0} f(t)$$

- Can also consider this as $\mathbf{E} \rightarrow \mathbf{E} + \eta \mathbf{J}_{\text{rf}}$, where $\eta \mathbf{J}_{\text{rf}}$ is an emf-per-unit-length inducing magnetic field in the plasma.

- Variable parameters $\lambda_{\text{rf}}, R_{\text{rf}}, Z_{\text{rf}}, w_{\text{rf}}$ (amplitude, location, spatial width), along with time dependence $f(t)$; begin with cylindrical equilibrium.
- Source *not* averaged over flux surface

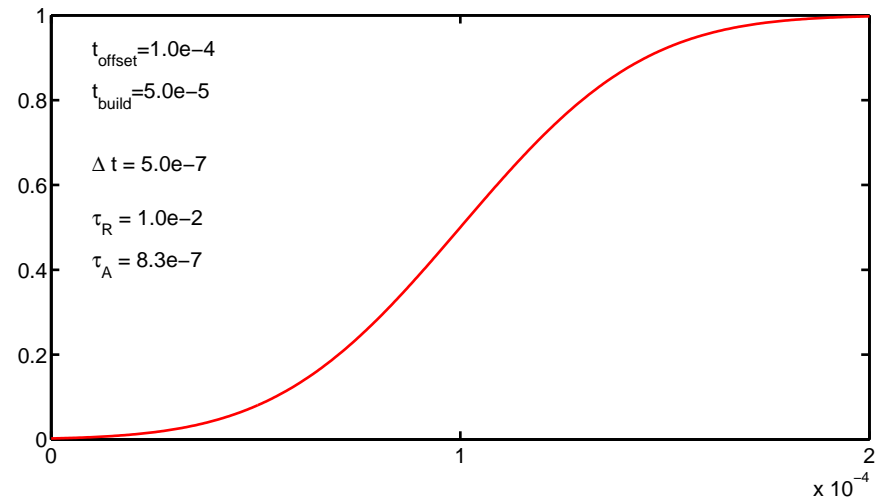


Ad hoc RF current can be slowly ramped up in time

- Specify the time dependence $f(t)$:

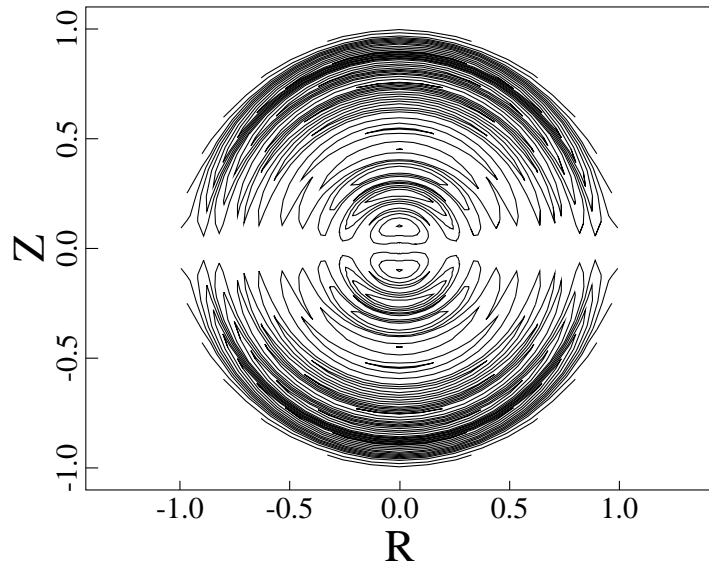
$$f(t) = \left[\frac{1}{2} + \frac{1}{2} \operatorname{erf} \left(\frac{t - t_{\text{offset}}}{t_{\text{build}}} \right) \right]$$

- Build up current on timescale t_{build} to suppress transient Alfvén waves (t_{build} short compared to resistive diffusion timescale τ_R)
- Initial perturbation has amplitude comparable to random current fluctuations; $\lambda_{\text{rf}} = 8.0 \times 10^{-4}$, $w_{\text{rf}} = 0.1$, $R_{\text{rf}} = 0.5$, $Z_{\text{rf}} = 0$



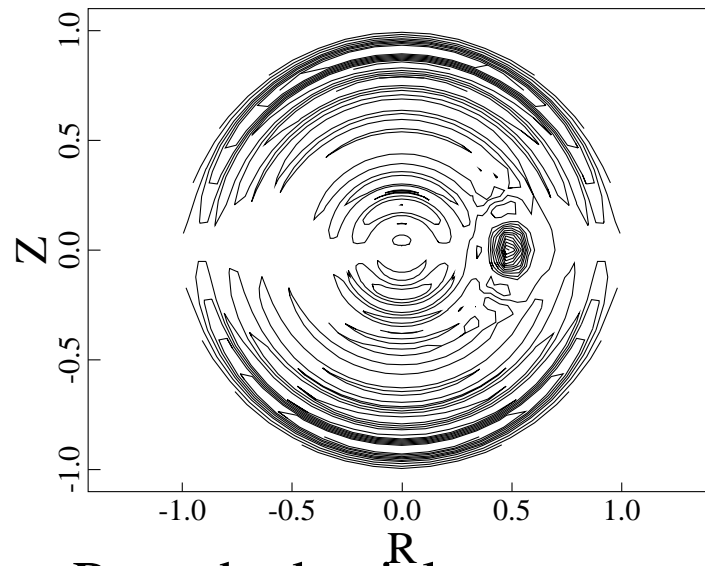
Current equilibrates over a flux surface in a few Alfvén times

Perturbed axial current



$t = 0$

Perturbed axial current

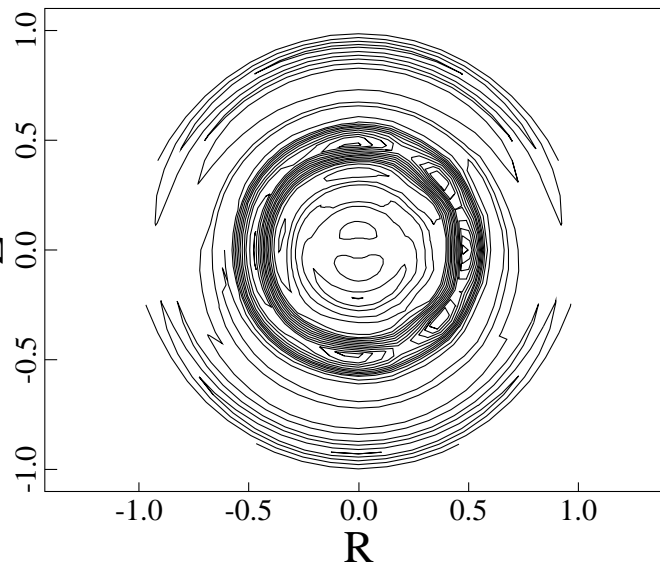


Perturbed axial current

$t = \Delta t$

- $\Delta t \approx 0.6\tau_A$, after only a few Alfvén times, current is spread across entire flux surface

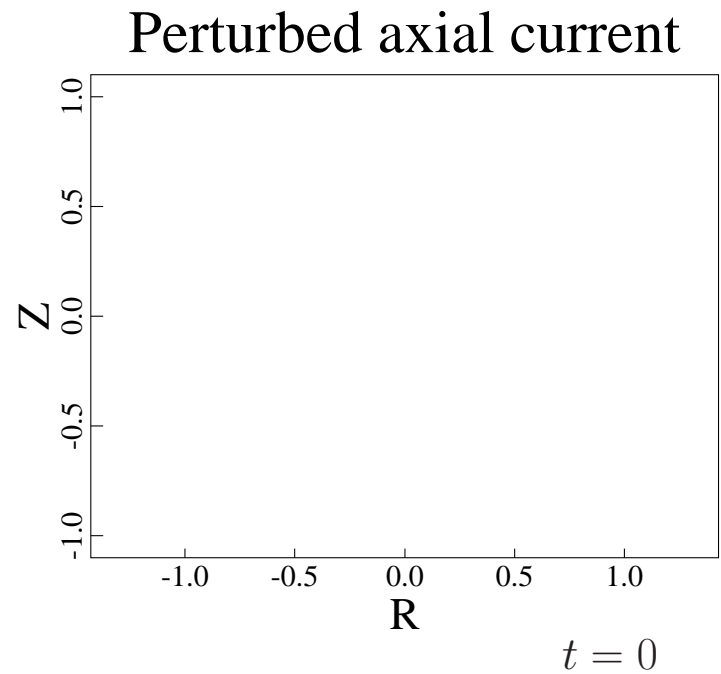
- Initial condition induces small-amplitude Alfvén waves (not related to RF source)



$t = 20\Delta t$

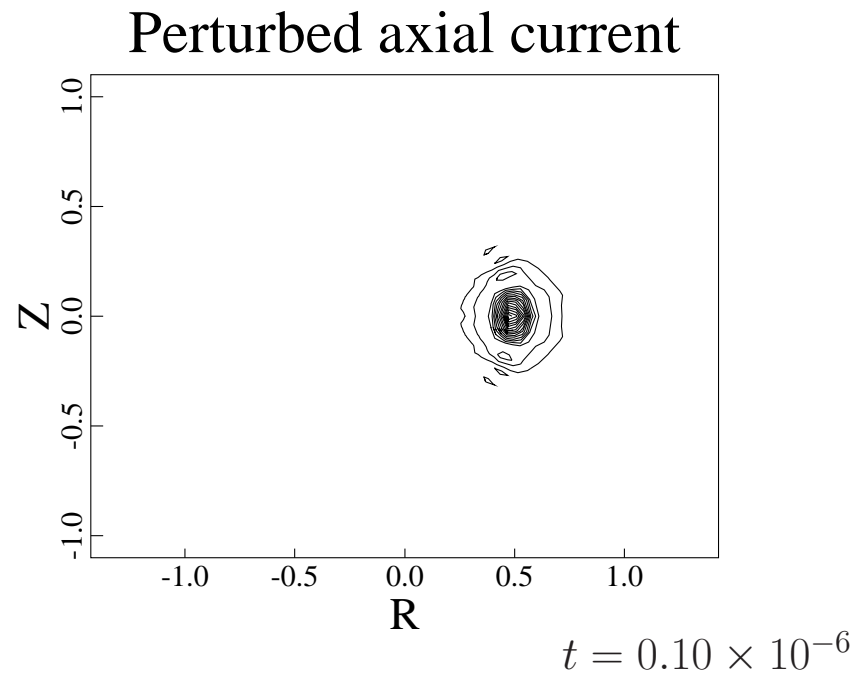
Current equilibration: Alfvén timescale

- Examine this process more closely by reducing the timestep and increasing λ_{rf}
- $\lambda_{\text{rf}} = 16.0, \Delta t = 1.0 \times 10^{-7}, \tau_A = 8.3 \times 10^{-7}$:



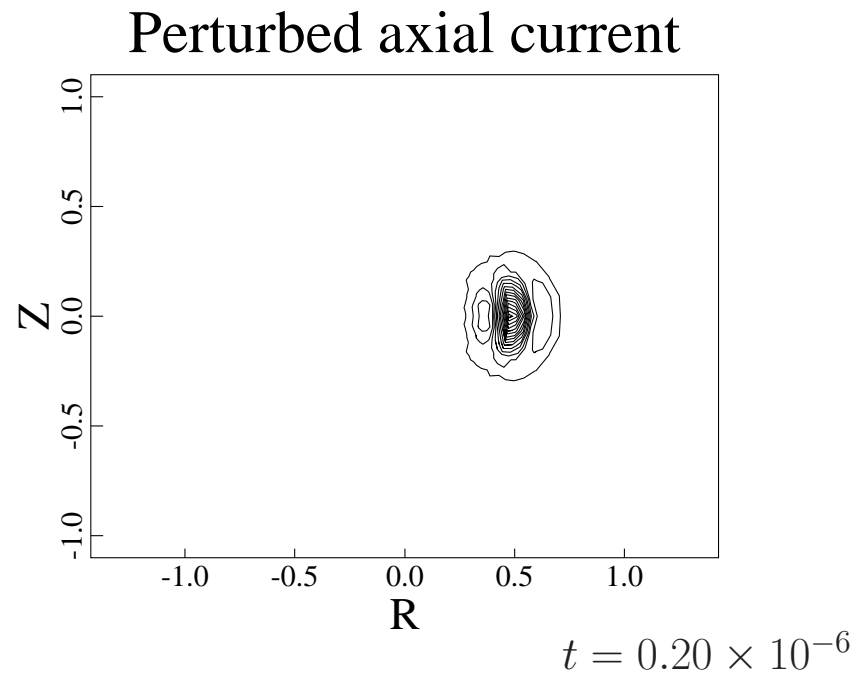
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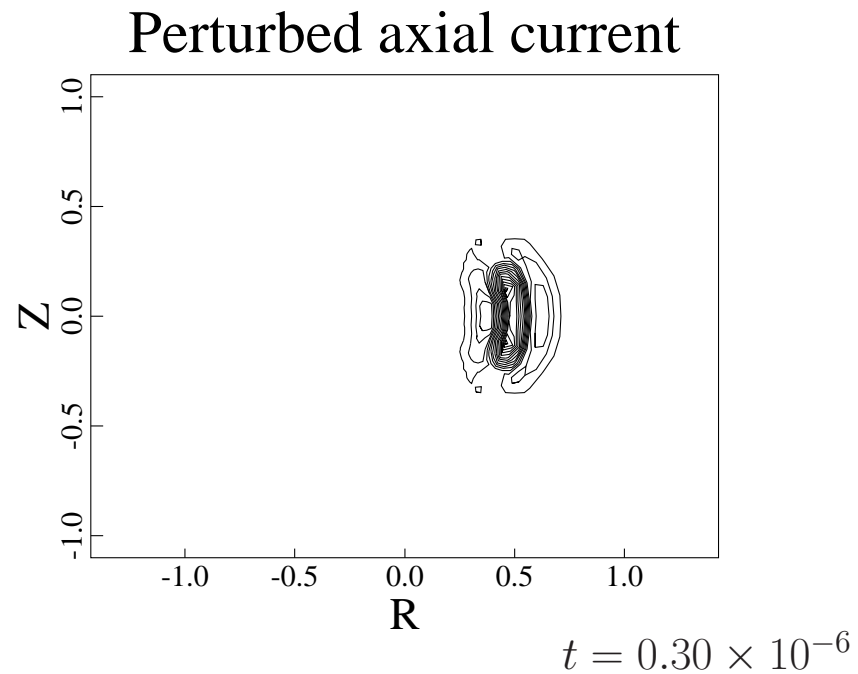
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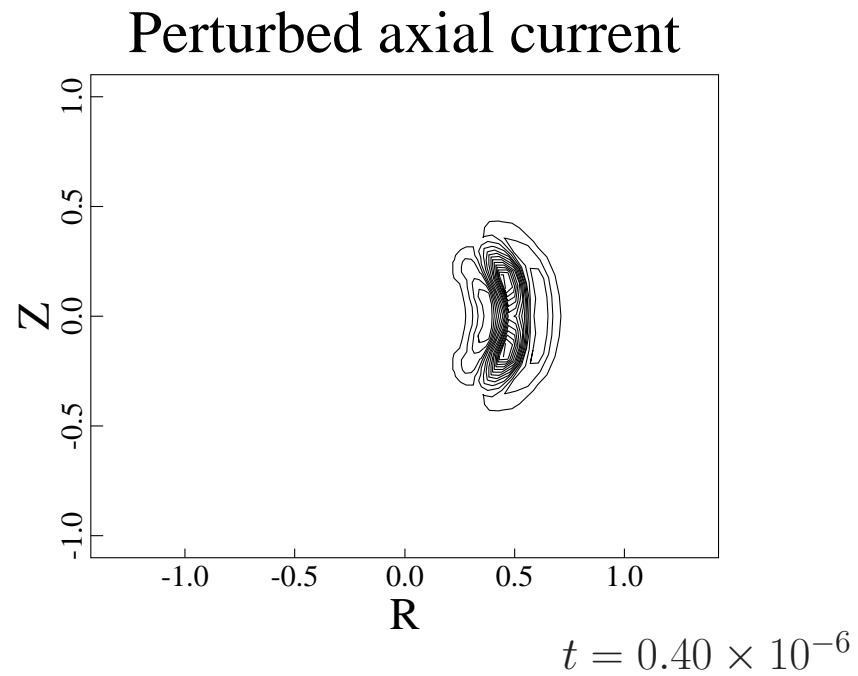
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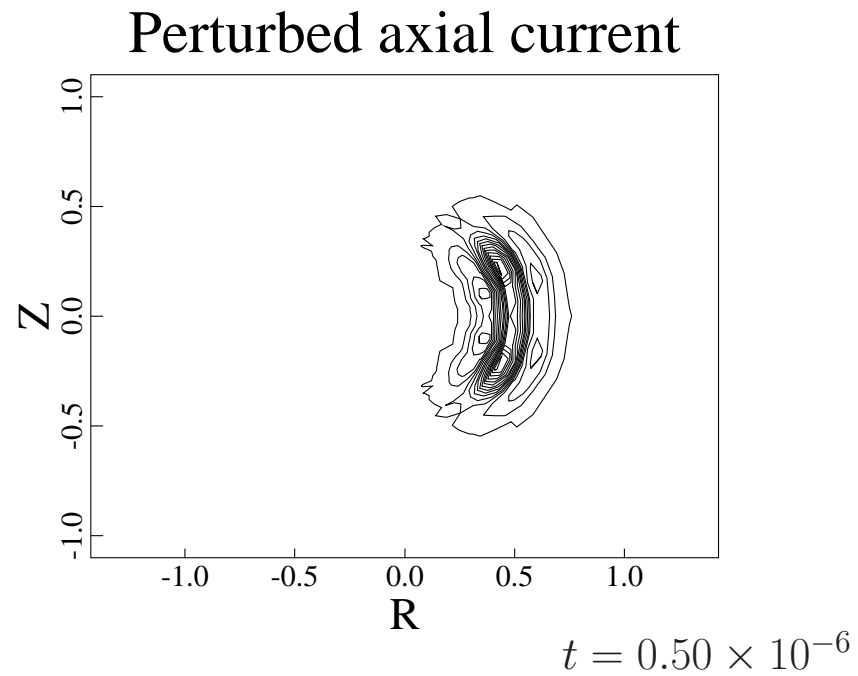
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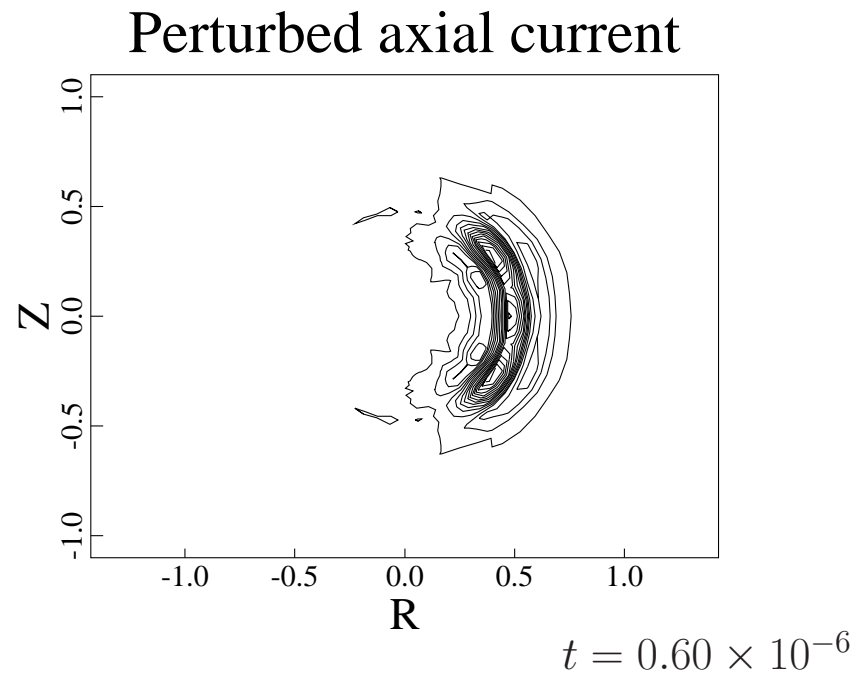
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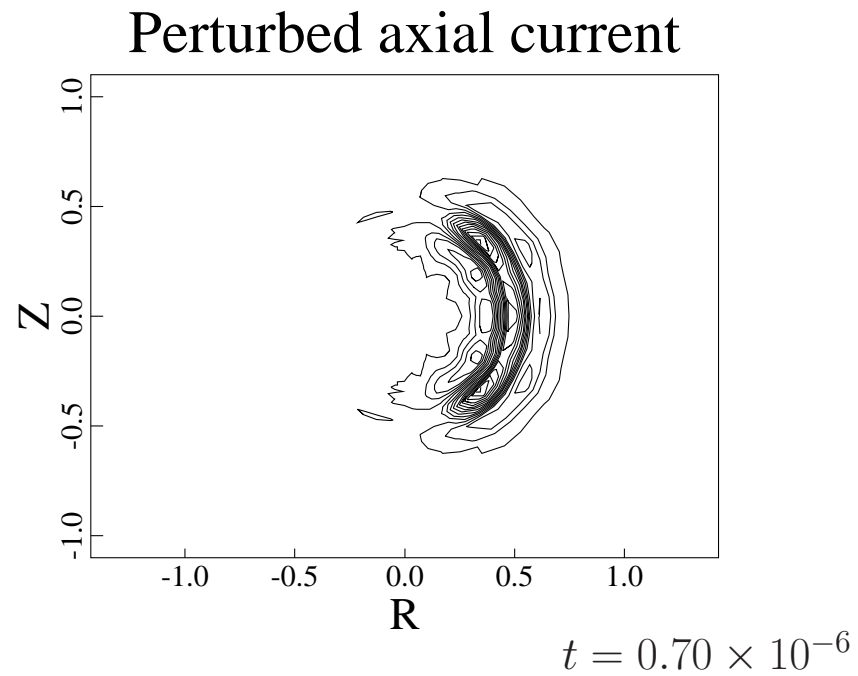
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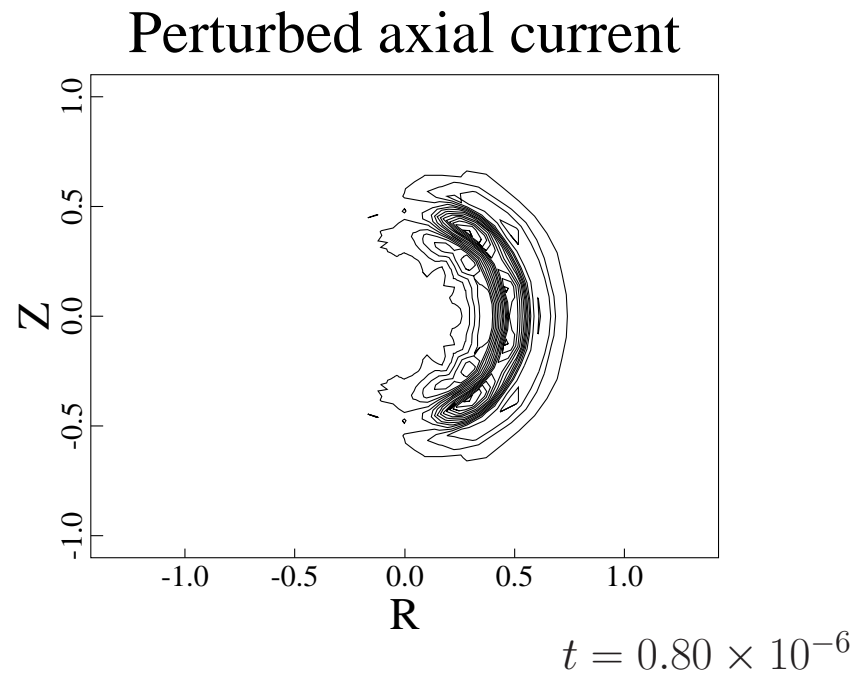
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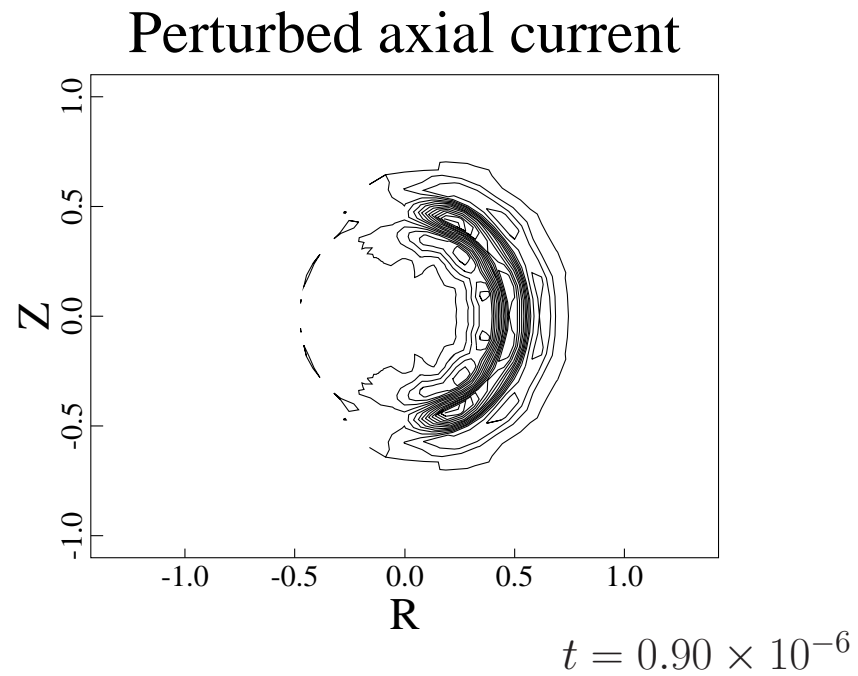
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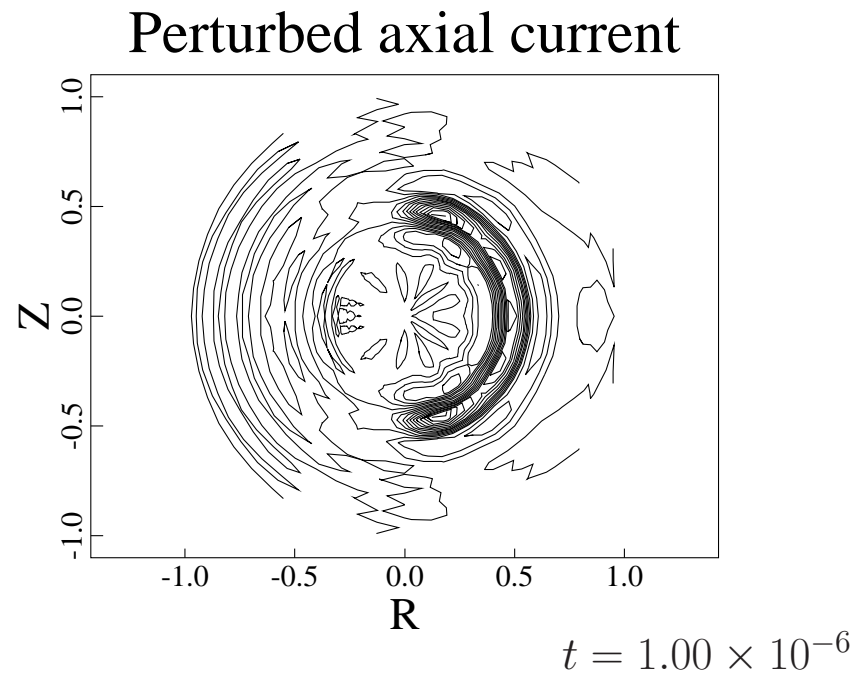
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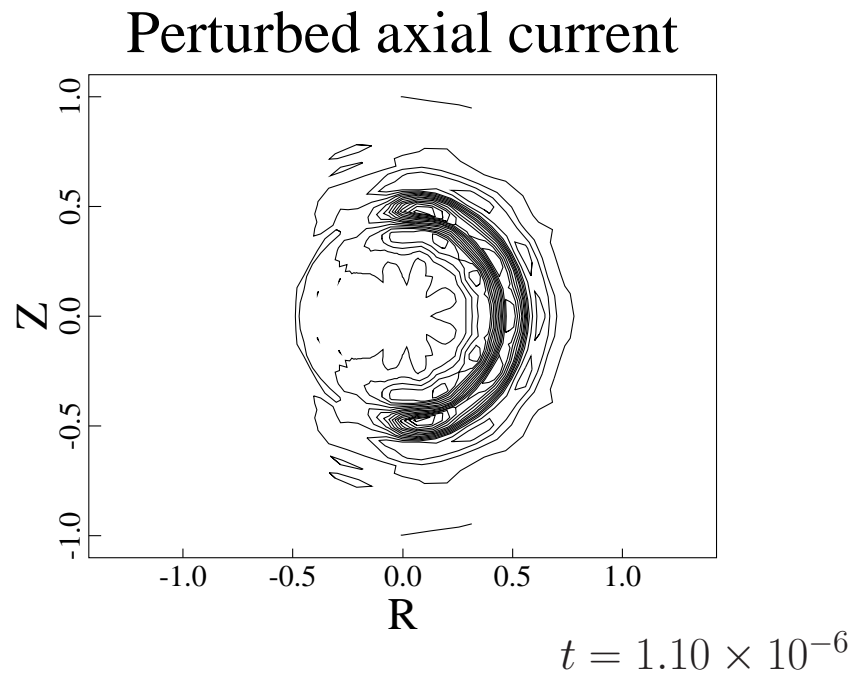
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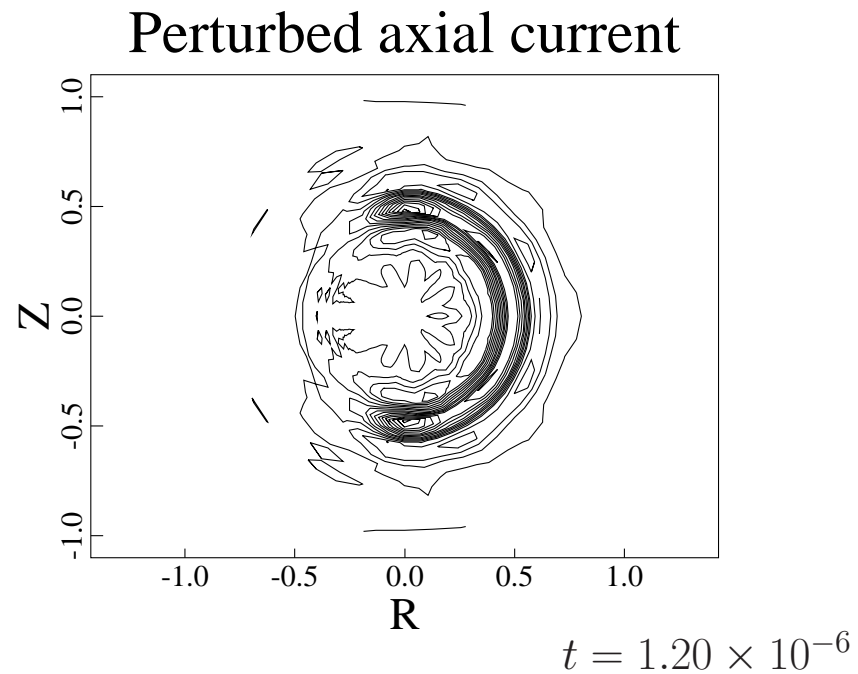
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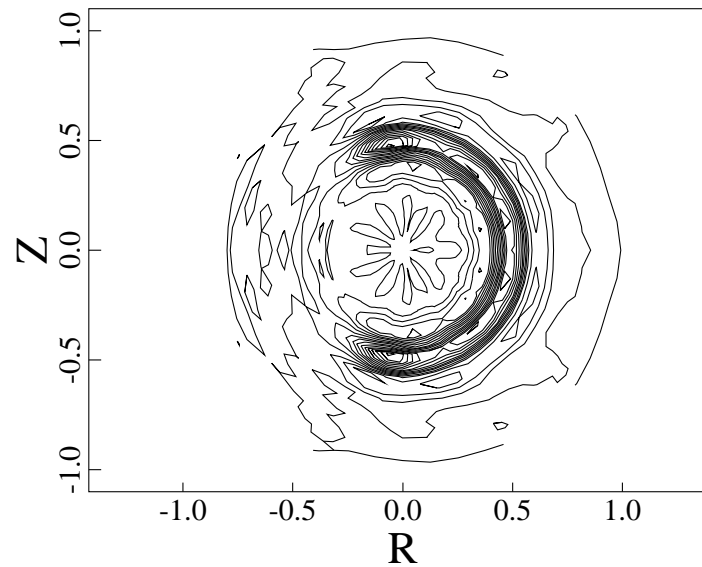
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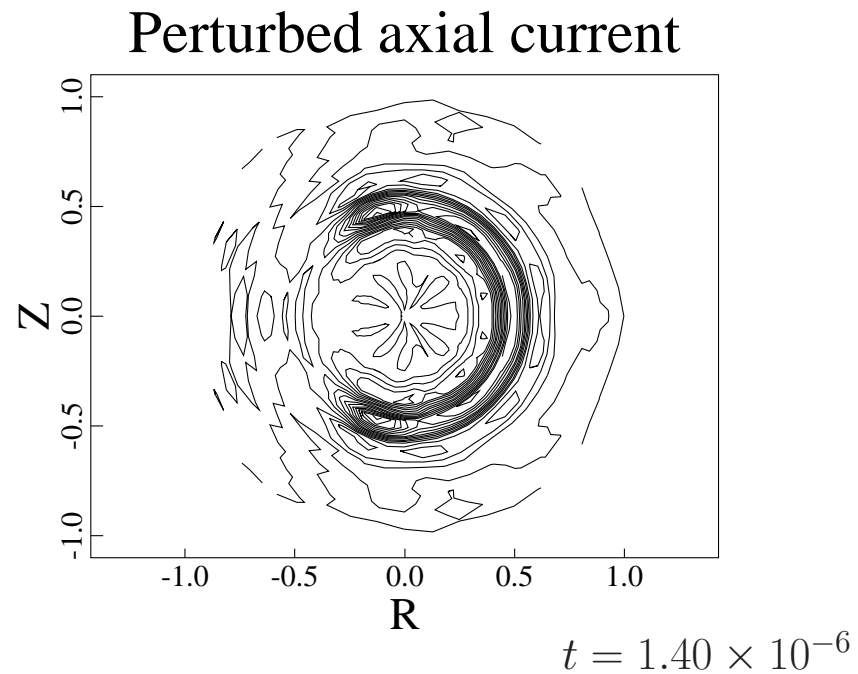
Perturbed axial current



$$t = 1.30 \times 10^{-6}$$

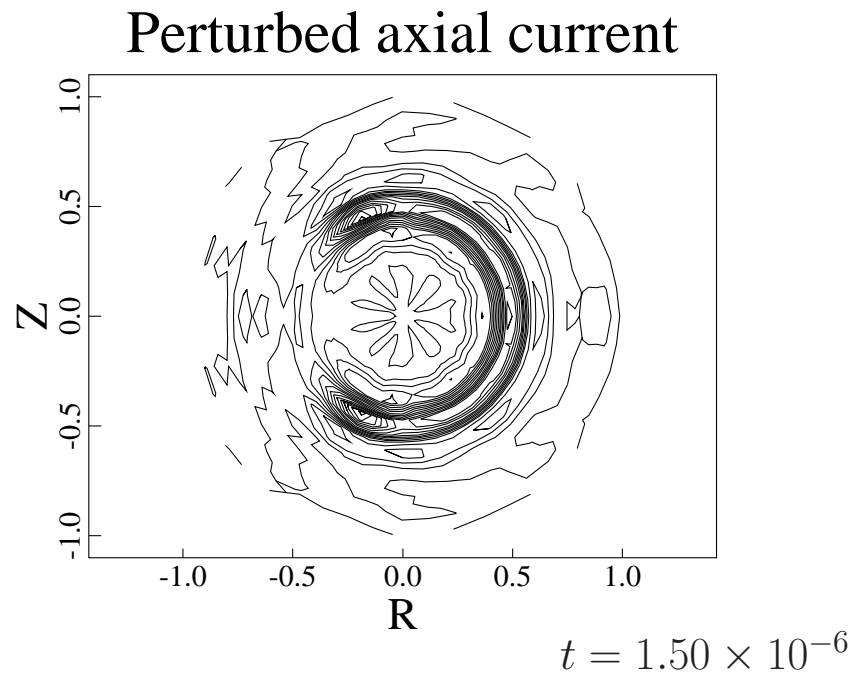
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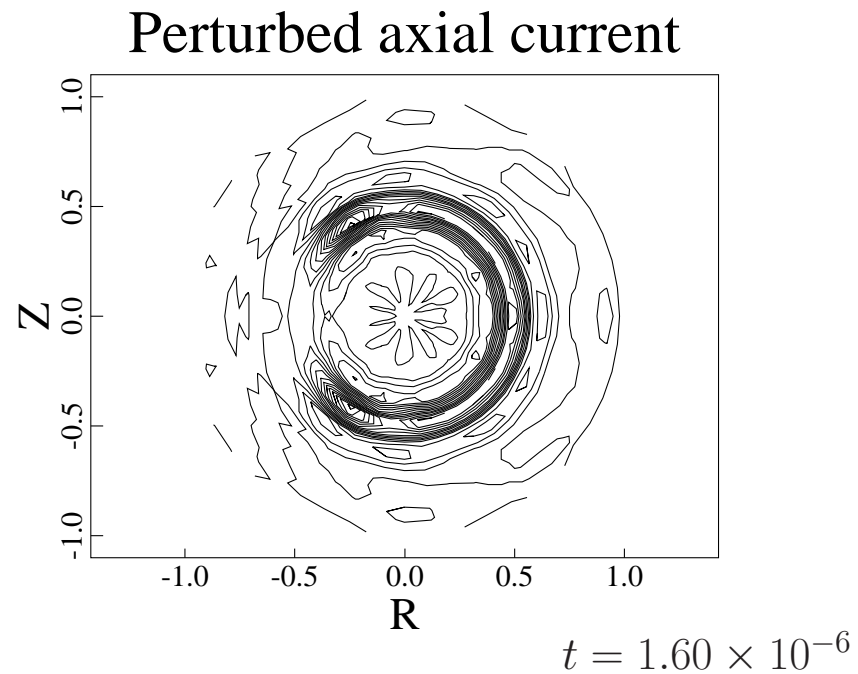
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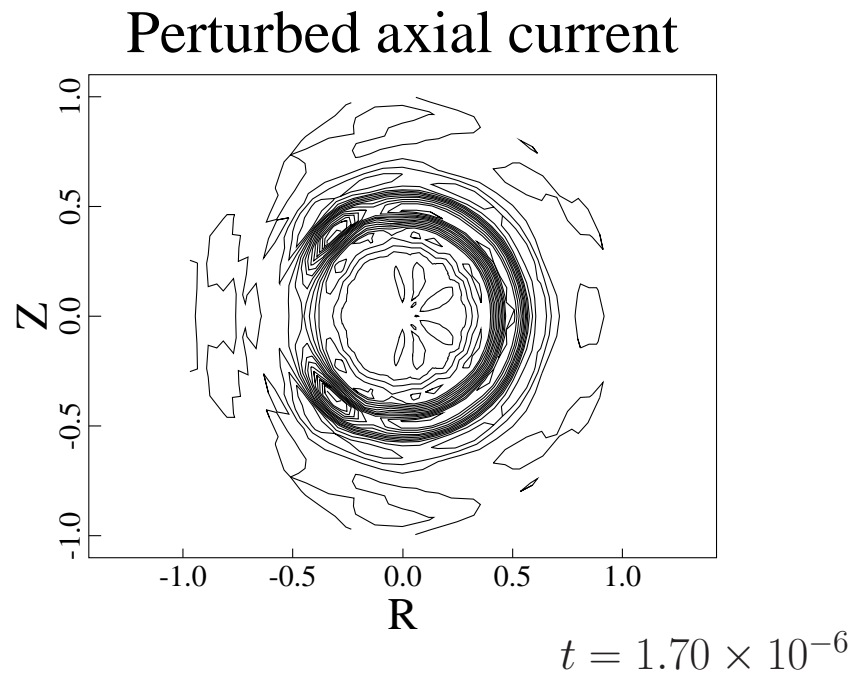
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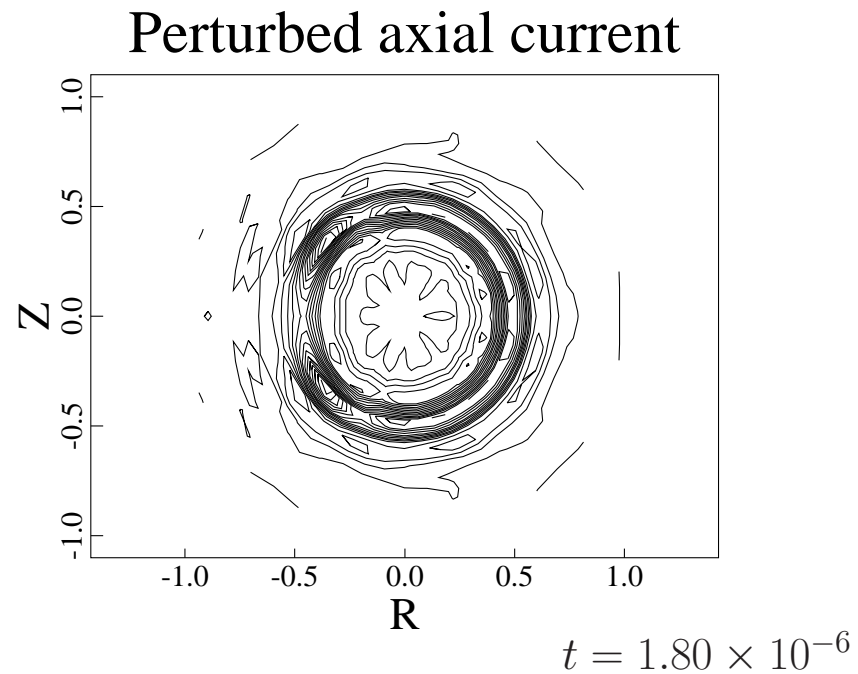
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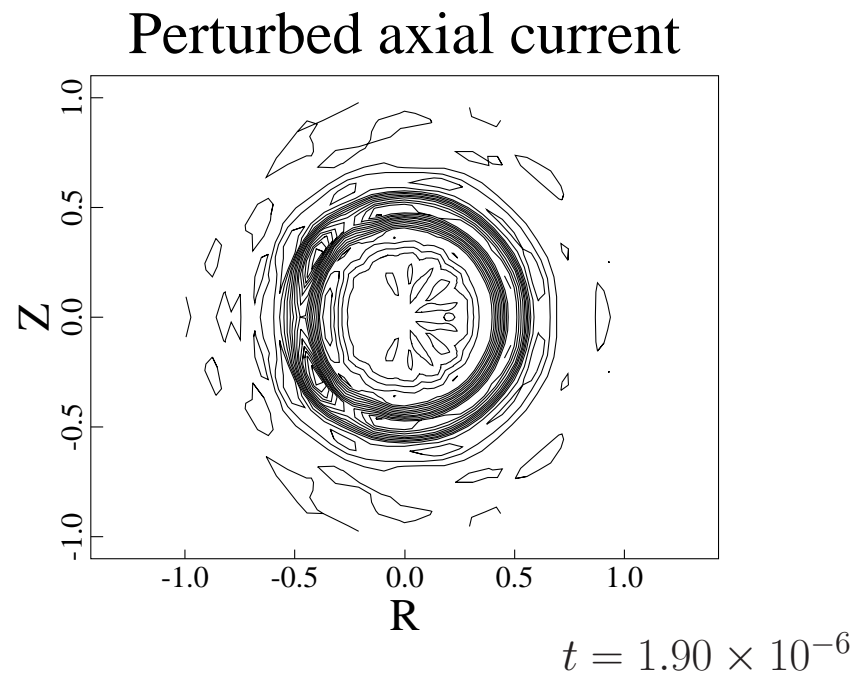
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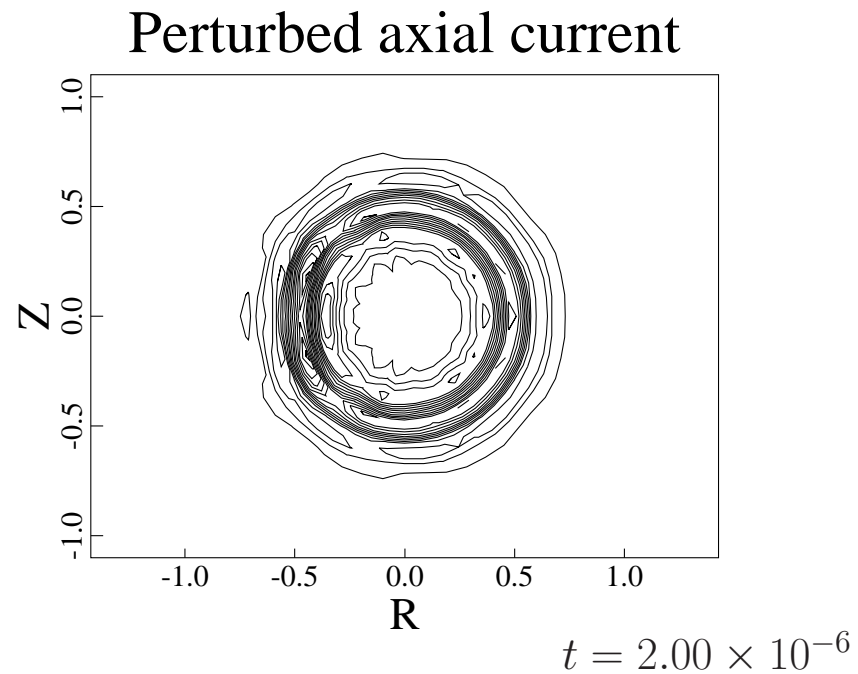
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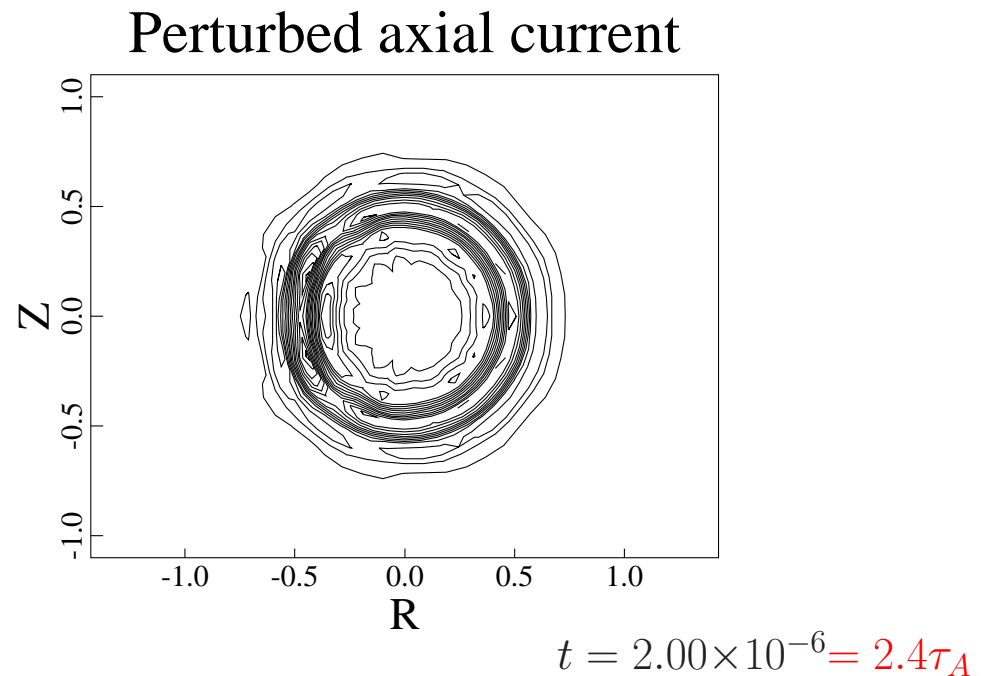
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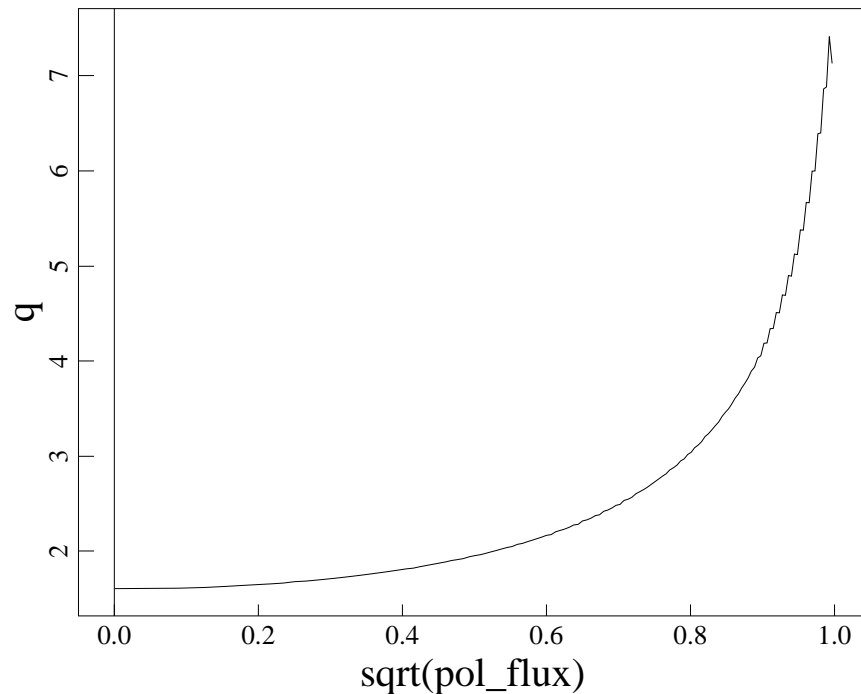


- Current equilibrates over a flux surface on the Alfvén timescale.

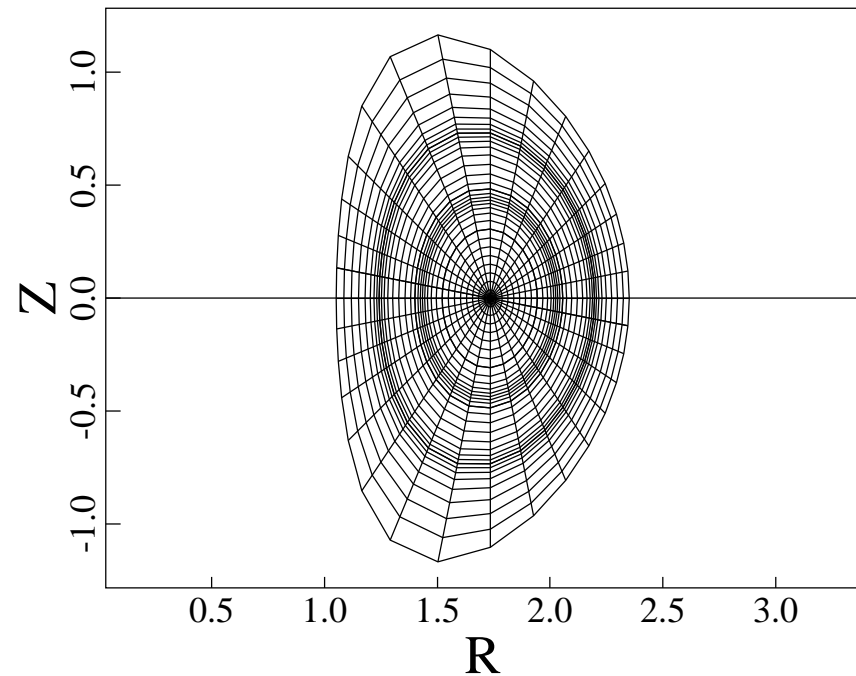
Test this effect in realistic (DIII-D-like) geometry

- Grid packing used to resolve rational surfaces ($q = 2, q = 3$) more accurately

q vs sqrt(pol_flux)



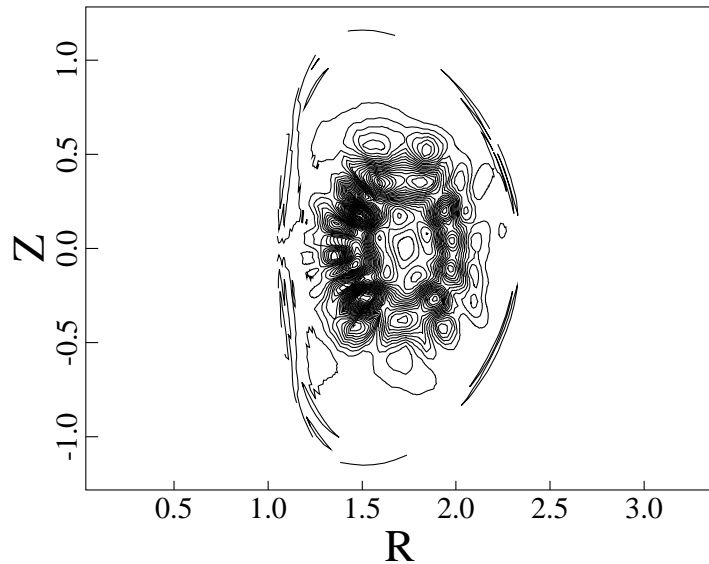
Finite Element Mesh



- Benchmark: DIII-D-like simulations of Sovinec *et al.* [J. Comp. Phys. **195**, 355 (2004)].
- Use axially symmetric \mathbf{J}_{rf} (a “ring” of current in the tokamak); examine effects on axially symmetric component of equilibrium ($n = 0$)

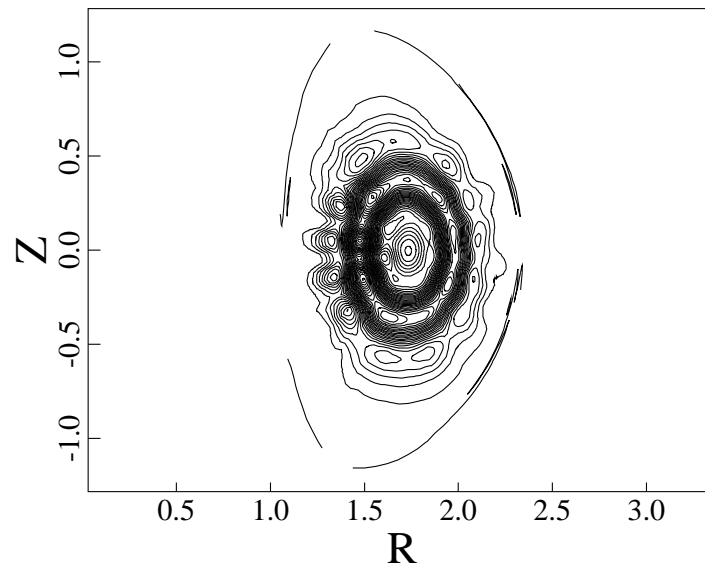
Current equilibration occurs in realistic geometries

Perturbed axial current



$$t = 600\Delta t = 1000\tau_A$$

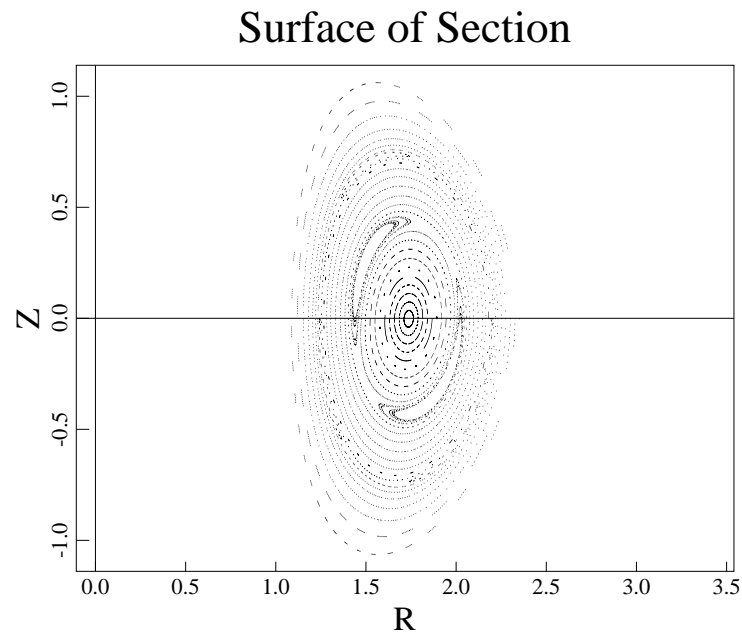
Perturbed axial current



$$t = 700\Delta t = 1170\tau_A$$

- Equilibration slower in realistic geometries; still fast compared to resistive diffusion time.
- $\lambda_{\text{rf}} = 1.0, Z_{\text{rf}} = 0, R_{\text{rf}} = 2.0, w_{\text{rf}} = 0.1$
- $t_{\text{offset}} = 0.02, t_{\text{build}} = 0.005, \Delta t = 1.0 \times 10^{-5}, \tau_R = 0.998, \tau_A = 6.0 \times 10^{-7}$
- $R/a \sim 3, S = 1.66 \times 10^6, Pr = 0.1$

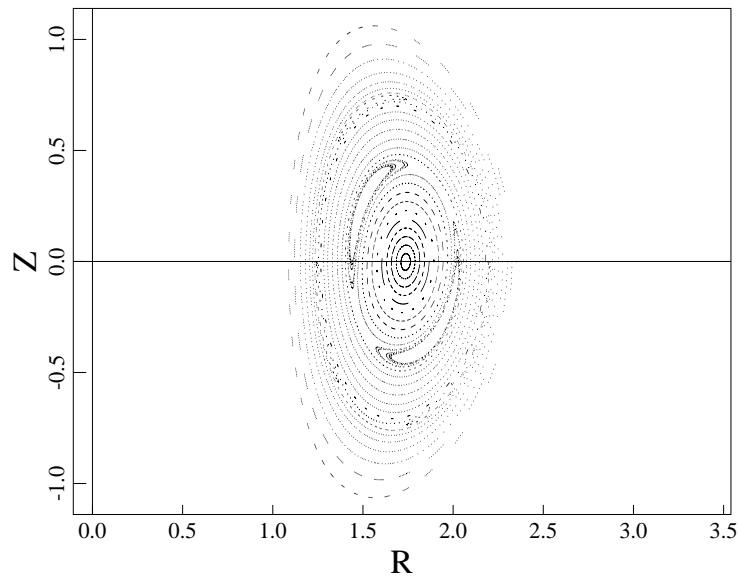
(2,1) magnetic islands grow/saturate in the absence of RF term



- Dominant nonlinear effect is (2,1) magnetic island
- Smaller island chain at $q = 3$ surface
- What happens when RF term is added?

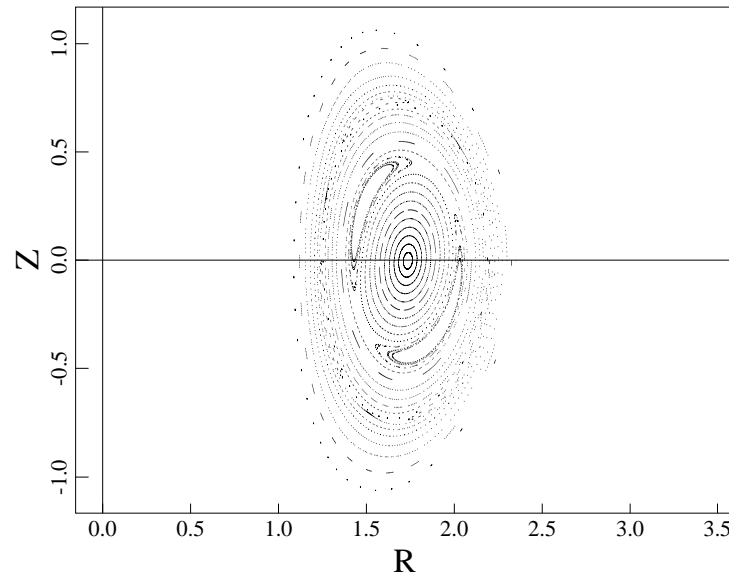
RF effects influence island size ... issues remain

Surface of Section



No RF

Surface of Section



RF

- Islands appear to be increasing slightly in size (wrong direction)
- Some issues with numerical stability of RF runs; working with Dalton Schnack and Carl Sovinec to address this
- Next steps: Investigation of effect of ad hoc current location/breadth on island size; non-axisymmetric sources, etc.